

# Electron beam experiments at FAST

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# Outline

## 1 FAST electron photoinjector

Facility and beamline overview

1.3 GHz SRF cavity transfer map measurement

## 2 Microlens array laser shaper

Laser transverse shaping, emittance reduction

Multi-beam generation and applications

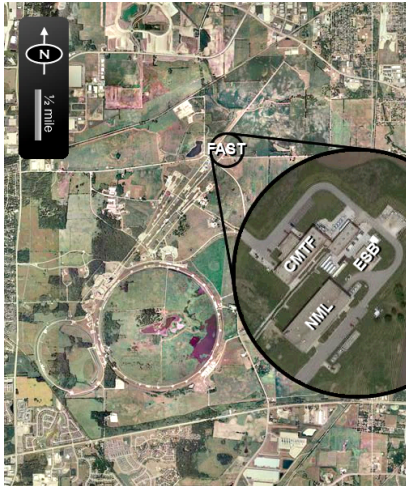
## 3 CAM and flat beam generation

CAM beams formation

Flat beam generation and emittance measurements

## 4 Longitudinal space-charge amplifier

# FAST introduction

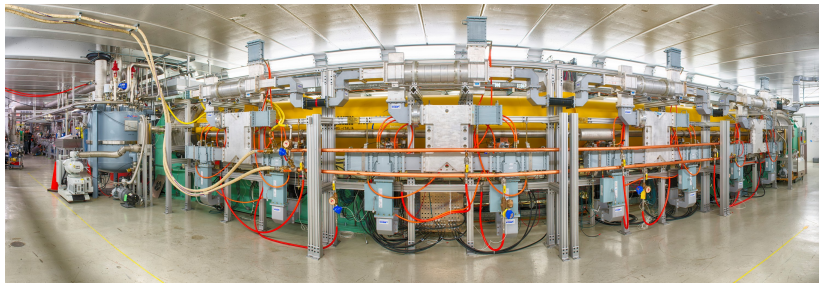


## Fermilab Accelerator Science and Technology - FAST

- 300 MeV electrons
- Linac + Ring
- End of construction - late 2018

<http://fast.fnal.gov/>

# FAST injector + IOTA ring

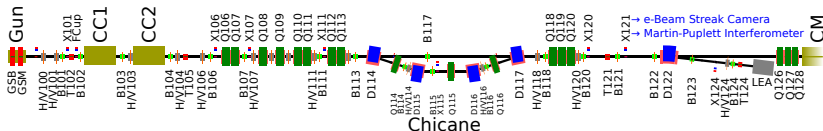


ILC-type cryomodule - Picture is courtesy of FAST

- $\text{Cs}_2\text{:Te}$  photocathode + 1.3 GHz RF gun
- Two 1.3 GHz SRF capture cavities + cryomodule = 300 MeV
- Injection into IOTA ring (150 MeV) + high energy electron beam experiments (X-ray channeling, ICS, flat beams)



# Electron injector



2015 (20 MeV) → 2016 (52 MeV) → 2017 (301 MeV)  
**2018 Ring completion / Experimental program start**

- Charge range: 10 fC - 3.2 nC per pulse (up to 3000 pulses/s)
- Nominal bunch length: 5 ps (minimum: 2 ps)
- Magnetic chicane and skew-quadrupole adapter (RTFB)
- Includes interaction points for medium (50 MeV) and high (300 MeV) energies, multislits, goniometer, pyro, etc.

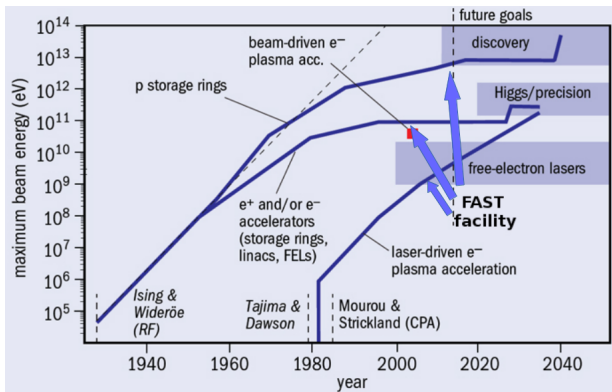
Detailed description: Antipov, S., *et al*, JINST, **12**, T03002 (2017)

# Electron beam parameters

Parameter	Value	Units
Emittance (norm.)	0.7	$\mu\text{m}$
Beam energy	50 - 300	MeV
Slice energy spread	<5	keV
Nominal charge	250	pC
Bunch length	5	ps
Beta-function (CC2 exit)	8	m
Dipole bending radius	0.958	m
Dipole length	0.301	m
Dipole angle	18	degrees
$R_{56}$	-0.18	m

Beam-based alignment: Romanov, A., arXiv:1703.09757  
[physics.acc-ph]

# Motivation for Research



Livingston plot - Image courtesy of CERN

How does electron beam research contribute to the field?

# Dissertation Impact

What we wanted to do:

*Electron beam transverse and longitudinal shaping in a photoinjector*

Why:

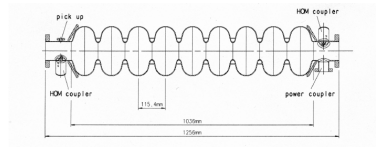
- ① Understand and improve beam dynamics at FAST
- ② Implement transverse laser shaper, improve emittance
- ③ Perform Round-to-Flat beam transform
- ④ Consider space charge amplifier at FAST

**FAST - Fermilab Accelerator Science and Technology facility**

<http://fast.fnal.gov/>

# 1.3 GHz SRF accelerating cavity

Beam dynamics of FAST low energy beamline defined by:

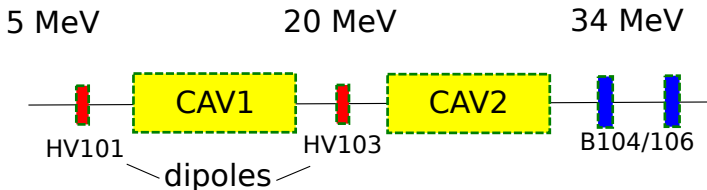


*Several proposed or operating accelerator facilities include TESLA-type cavities, such as FAST, ILC, LCLS-II, PIP2 and etc. to accelerate electron, proton or muon beams*

- 1 Experimentally verify Chambers-Serafini-Rosenzweig model
- 2 Attempt to characterize the effects of couplers

# Experimental setup (2016)

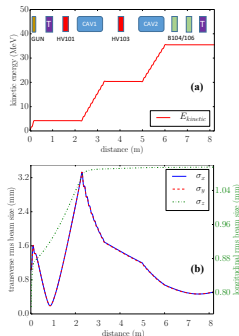
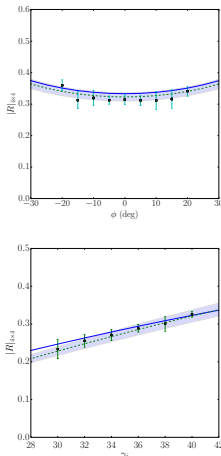
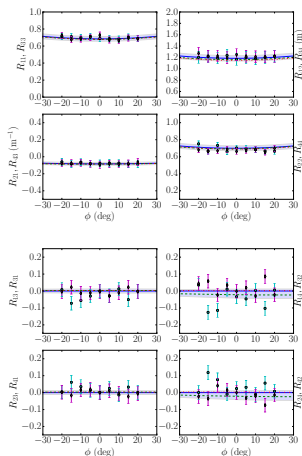
## *Schematics of the experiment*



## **Experiment details**

- ① Diagnostics/controls automatic (Piot, Halavanau (NAPAC16))
- ② Possible to vary injection energy (use CAV1, measure CAV2)
- ③ Reference orbit method;  $R$  inverted with least squares
- ④ Strong focusing in CAV1 (alters beam quality)
- ⑤ Instrumentation (BPM jitter  $< 80$   $\mu\text{m}$ ), laser

# Results



Each data point is inferred from 80 reference orbits

(left) transfer matrix  $R$  elements; (middle) determinant  $R_{4 \times 4}$  as a function of phase ( $\phi$ ) and injected  $\gamma_i$ ; (right) beam dynamics in low energy section.

# Cavity transport summary

## Conclusions:

- Chambers' model is accurate at FAST energies ( $>34$  MeV)
- HOM effect - phase dependent parametric dipole kick

## Outcomes:

- Beam-based alignment can be done via minimization procedure (experimentally confirmed for CG/BFGS-methods)
- Better understanding of low energy round beam dynamics, helps with flat beam
- Improved analytical linear model of linac (used for 300 MeV commissioning)
- Tools (pyACL, beam-based alignment)

*Halavanau, A., Phys. Rev. Accel. Beams* **20**, 040102 (2017)



# Emittance studies

Nominal FAST electron beam norm. emittance  $\epsilon = 0.7\mu\text{m}$  at  
comissioning charge of  $Q=250\text{ pC}$  (small laser spot +  
optimization)

Available measurement techniques:

- 1 Quadrupole scan (automatic)
- 2 Horizontal/vertical multislits
- 3 Possible to install pepper-pot

FAST electron beam norm. emittance at fully opened laser iris  
 $\epsilon = 1.9\mu\text{m}$  ( $\sigma \approx 1\text{mm}$ )

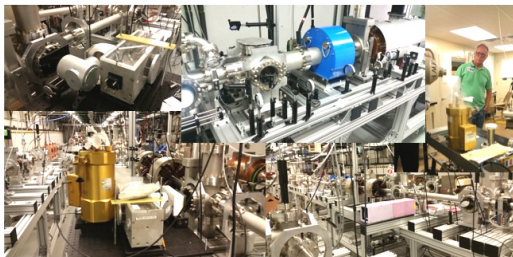
# How to reduce emittance by x2?

$$\epsilon = \langle \sigma_{\perp}^2 \rangle^{1/2} \langle \Delta\theta_{\perp}^2 \rangle^{1/2}, \quad \Delta\theta_{\perp}^2 = \mathcal{F}(T_{\text{eff}} + F_i + F_{SC})$$

$F_{SC}$  can be linearized in the case of transverse uniform distribution

**Laser can be homogenized by Microlens Arrays (MLAs)**

*Inspiration: bumpy ceiling light cover*

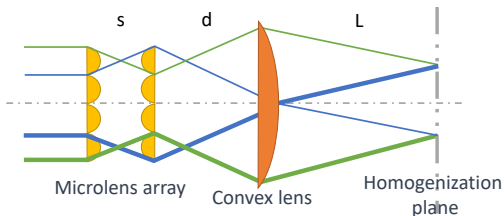
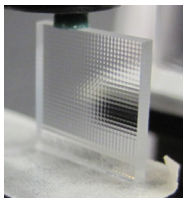


*72 MeV photoinjector + EEX beamline. Proof-of-principle MLA shaping experiment, emittance reduction by factor of 2, commissioned and used for experiments at AWA*

# Microlens arrays (MLAs)

In photocathodes the achievable electron beam parameters are controlled by the laser used to trigger the photoemission.

*Microlens arrays are fly-eye type light condensers*

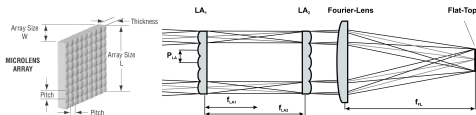


- 1 Produce uniform laser image in the focal plane of the mixing lens
- 2 Produce transversely modulated laser beams

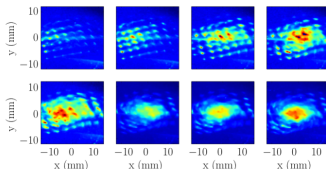
# Microlens array setup

- 1 Homogenized/Patterned beam can be imaged (4 lens solution)
- 2 Can produce high intensity beams
- 3 Hexagonal lattice for best homogenization

## Microlens array (MLA) setup at AWA



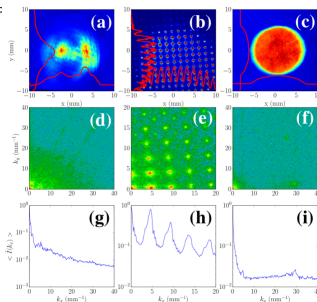
**Multi  
beam at  
50 MeV**



Halavanau, et. al, Phys.Rev.Accel.Beams 20 (2017), 103404

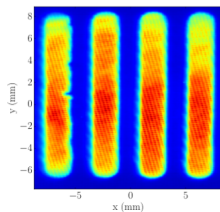
## AWA UV laser

Regular beam    Multi beam    Uniform beam

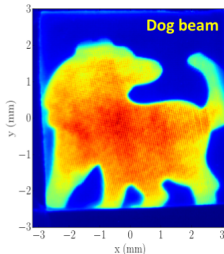
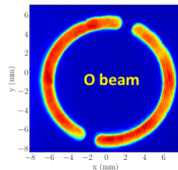
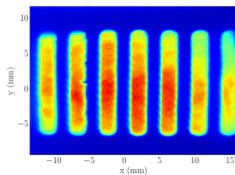


# MLA laser shaper

*Arbitrary laser transverse profile: homogenizer + mask*



Stacked beams

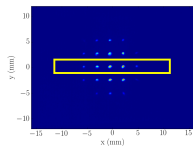


MLAs were mounted on a rotation stage; pinhole

# Emittance exchange setup

- 1 Use MLA to produce multi-beams
- 2 Send multi-beam through EEX and generate bunch trains
- 3 Use MLA rotation for bunch train tuning

## Experiment schematics: (MLA + EEX)

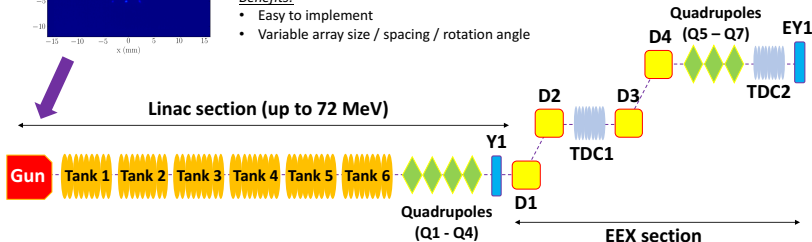


**Goal:** Combine MLA multi-beams with EEX to generate bunch trains

- MLA setup can shape laser spot to create multi-beams
- Additionally, a line of laser dots can be selected

**Benefits:**

- Easy to implement
- Variable array size / spacing / rotation angle



In progress, reported at IPAC18

# Microlens array summary

- ① Generated homogenized and patterned beams with a single setup (elegant and simple)
- ② Comissioned and used routinely at AWA
- ③ Application in photocathode quantum efficiency measurement (**NEW, in progress**)
- ④ Application in bunch train generation (**NEW, in progress**)
- ⑤ Implementation at FAST underway
- ⑥ Interest of SLAC, UCLA, LBNL, PITZ and many others

# Why magnetized beams?

## **Canonical angular momentum (CAM) dominated beams:**

- ① Conventional application - electron cooling (Derbenev, Ya., UM-HE-98-04-A)
- ② Emittance partitioning via flat beams (interest of AWA group)
- ③ Flat beams in plasma acceleration (interest of UCLA/AWA)
- ④ Flat beams in DLWA (interest of PEGASUS facility)
- ⑤ Suppressing microbunching instabilities in IOTA (collaboration with R. Li, JLab)
- ⑥ Several possible radiation experiments (dielectric structures, microundulators, channeling, etc.) can be done at FAST

**CAM beams production at FAST is a stepping stone**



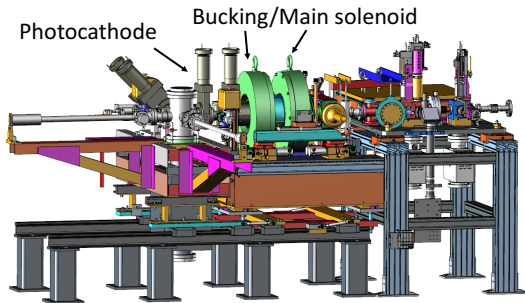
# Busch's theorem

*Total canonical angular momentum  
of a charged particle in symmetric magnetic field is conserved*

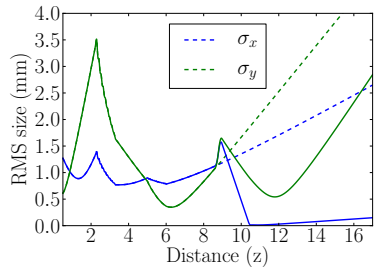
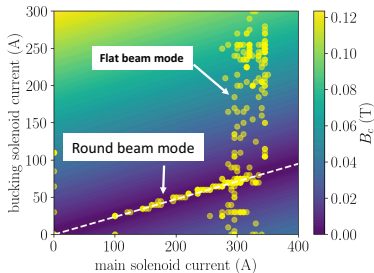
$$L = \gamma m r^2 \dot{\theta} + \frac{1}{2} e B_z(z) r^2 \quad \mathcal{L} = L/2p_z$$

Eigenemittances:

$$\epsilon_{\pm} = \sqrt{\epsilon_u^2 + \mathcal{L}^2} \pm \mathcal{L} \rightarrow \epsilon_+ \approx 2\mathcal{L}; \quad \epsilon_- \approx \frac{\epsilon_u^2}{2\mathcal{L}}$$



# CAM and flat beam dynamics



- 1 Two gun solenoids must ensure full transmission  $\rightarrow$  can't wire them opposite (in that case max  $B_z=0.2$  Tesla)
- 2 In our experiment  $B_z=0.07$  Tesla was selected (after solenoid optimizations)
- 3 Dash/solid lines represent magnetized/flat beam RMS size

# Emittance ratio

Eigenemittances:

$$\epsilon_- \equiv -\sqrt{\epsilon_0^2 + \mathcal{L}^2 - 2\mathcal{L}\epsilon_0} = -\sqrt{(\epsilon_0 - \mathcal{L})^2} = \mathcal{L} - \sqrt{\mathcal{L}^2 - \epsilon_4^2} \approx \frac{\epsilon_4^2}{2\mathcal{L}}$$

$$\epsilon_+ \equiv \sqrt{\epsilon_0^2 + \mathcal{L}^2 + 2\mathcal{L}\epsilon_0} = \sqrt{(\epsilon_0 + \mathcal{L})^2} = \mathcal{L} + \sqrt{\mathcal{L}^2 + \epsilon_4^2} \approx 2\mathcal{L}$$

Emittance ratio or “flatness”:

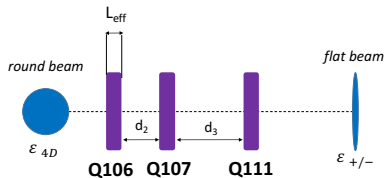
$$\frac{\epsilon_+}{\epsilon_-} = \frac{4\mathcal{L}^2}{\epsilon_u^2} = \frac{1}{p_z^2} e^2 B_{0z}^2 \frac{\sigma_0^2}{\sigma_0'^2}$$

Example calculation:  $\sigma_+ = \sqrt{\beta_{x,y}\epsilon_+} \rightarrow \epsilon_4 = 2 \mu\text{m} \rightarrow \epsilon_+ = 40 \mu\text{m}$ ,  
 $\epsilon_- = 0.1 \mu\text{m} \rightarrow \beta_{x,y} = 8\text{m}$ ,  $\sigma_+ = 1.8\text{mm}$  and  $\sigma_- = 0.09\text{mm}$

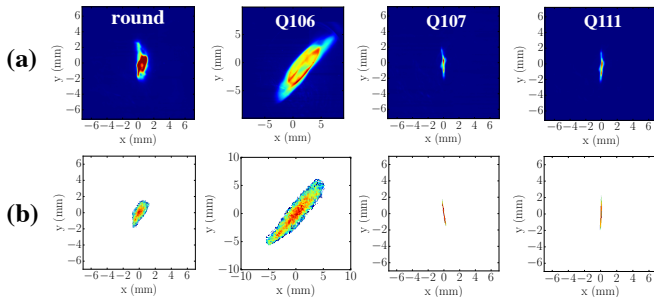
*Burov, A., Phys. Rev. E **66**, 016503 (2002)*

*Kim, KJ., PRSTAB, **6**, 104002 (2003).*

# Round-to-flat transformation



- Q106, Q107, Q111 - skew-quadrupoles
- (a) - Experimental, (b) - Simulations in Impact-T



Good agreement - good model!

# RTFB solutions (thin lens)

FAST quadrupoles:  $K = (10.135 \times 40 I_q) / (1.8205 \times p [\text{MeV}/c])$ ,  
 $L_{\text{eff}} = 17 \text{ cm}$

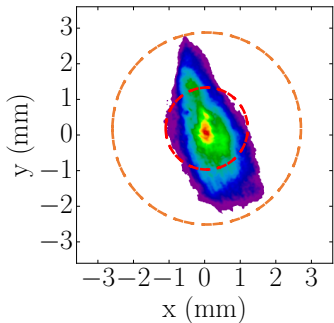
$$q_1 = \pm \sqrt{\frac{-d_2(d_T s_{21} + s_{11}) + d_T s_{22} + s_{12}}{d_2 d_T s_{12}}},$$

$$q_2 = \frac{(d_2 + d_3)(q_1 - s_{21}) - s_{11}}{d_3(d_2 q_1 s_{11} - 1)},$$

$$q_3 = \frac{d_2(q_2 - q_1 q_2 s_{12}) - s_{22}}{d_2(d_3 q_2 s_{22} + q_1 s_{12} - 1) + d_3(s_{12}(q_1 + q_2) - 1)}$$

*Numerical optimization can be used for correcting  $(q_1, q_2, q_3)$  for chromaticities and other second order effects*

# What if beam is not round?



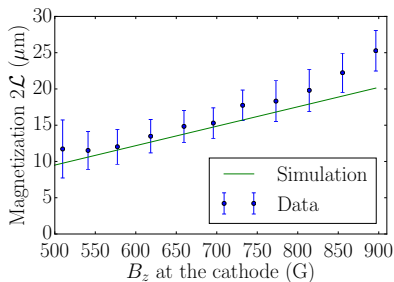
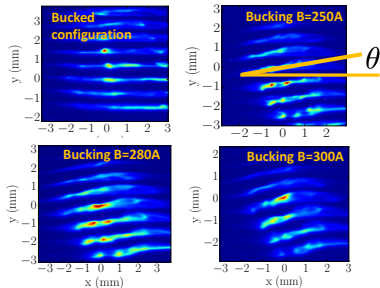
FAST laser cathode distribution

$$\sigma_x = 520\mu\text{m}, \sigma_y = 920\mu\text{m}$$

**First** flat beam with *asymmetric* laser!

- ① Assume very low charge (20 pC)  $\rightarrow$  no space charge. RTFB solutions do not depend on  $\mathcal{L}$ . White areas will be not present in the final phase space.
- ② When space charge is included, the problem requires 4 skew quadrupoles in RTFB setup
- ③ FAST Run 2017 used 3 magnets, will add additional in the future

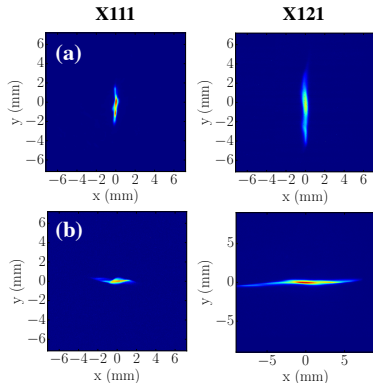
# CAM measurement with slits



- $\langle L \rangle = 2p_z \frac{\sigma_1^2 M \sin \theta}{D}$ , where  $p_z$  is momentum,  $D$  is the drift length,  $\sigma_1 = (n - 1) * d/5$ ,  $M = \sigma_2/\sigma_1$  - magnification factor
- First used in Fermilab A0 flat beam experiment (Sun, et. al.)
- Similar idea with multi-beam generated by MLAs (Halavanau, et.al, Phys. Rev. AB, **20**, 10, 103404, (2017))

# Vertical/Horizontal flat beams

$Q = 20 \text{ pC}$ ,  $B_z = 730 \text{ Gauss}$



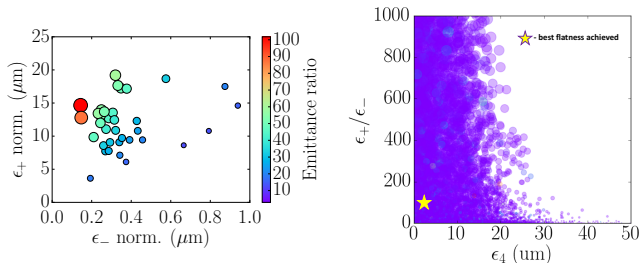
X111/X121 - screens 4 m apart  
downstream of RTFB

- Vertical flat beam  $\epsilon_- \rightarrow \epsilon_x$ , RTFB: + - +
- Horizontal flat beam  $\epsilon_- \rightarrow \epsilon_y$ , RTFB: - + -
- Beam-based optimizer: optimizing projections/ratio (not very efficient because  $\sigma = \sqrt{\beta\epsilon}$ )
- Emittances: (2 nm, 220 nm) geom., both hfb/vfb
- How to further optimize the emittance?



# FAST flat beam parameter space

(left) Experimental flat beam realizations at FAST. Size/color of circles defines aspect ratio. First automatic RTFB transformation!

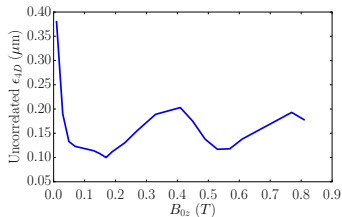
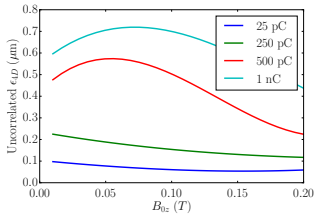


(right) 100,000 realizations of genetic optimization algorithm (MOGA). Optimizing flatness using: gun phase, gun gradient, CAV1/CAV2 parameters, spot size and solenoidal fields as variables (path to AI phase-space manipulation w/ Auralee Edelen).

# Further optimization

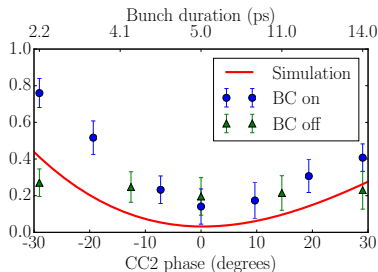
Idea by S. Nagaitsev:

**Can we compensate space-charge with strong field?**

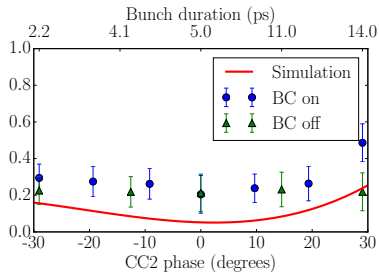


Preliminary cathode design considerations *in progress!*

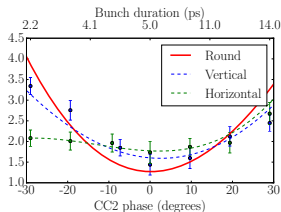
# First compressed flat beams!



Compressed vertical flat beam - significant emittance growth at maximum compression



Horizontal flat beam - small emittance in the same plane as chicane CSR, slight growth (Zhu, 2014)



- Horizontal flat beam emittance is largely unaffected by chicane CSR
- Total ( $\epsilon_x \epsilon_y$ ) preserved better

# Flat beam summary

- ① Generated CAM/flat beam from asymmetric laser (**NEW**)
- ② Automatic horiz./vert. flat beam transformation (**NEW**)
- ③ Lowest emittance  $0.1 \mu\text{m}$  (below thermal) (**NEW**)
- ④ Compressed flat beams, helps with beam transport (**NEW**)
- ⑤ AI phase-space manipulations (**NEW, in progress**)
- ⑥ Getting closer to ILC-type beam (**NEW, in progress**)
- ⑦ New comprehensive image analysis tool

## Future of flat beams at FAST:

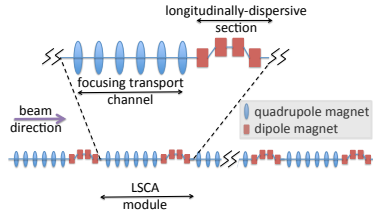
- ① High-charge flat beams (with J. Rosenzweig)
- ② Additional diagnostics → improve emittance ratio
- ③ Radiation generation at FAST (channeling, dielectric)

# Longitudinal space-charge amplifier

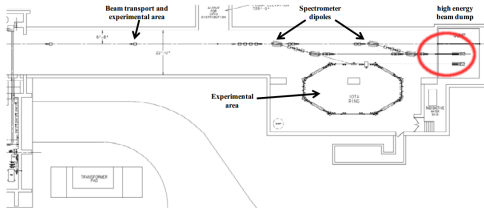
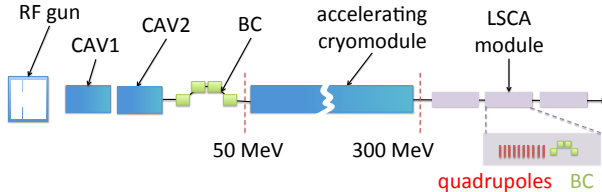
- Longitudinal space charge effects are responsible for unwanted energy modulations and emittance growth in FELs
- Can we take advantage of them?\*
- The technique was recently demonstrated in the optical domain\*\*

\*M. Dohlus, E. A. Schneidmiller, and M. V. Yurkov, *Phys. Rev. ST Accel. Beams*, **14**, 090702 (2011).

\*\*A. Marinelli, et al., *Phys. Rev. Lett.*, **110**, 264802 (2013).



# Possible location at FAST



Possible use of the FAST beamline before the high-energy adsorber area

# Space charge calculation

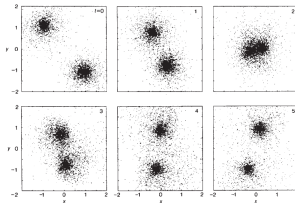
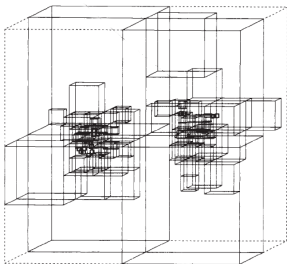
- Many numerical and analytical methods “reduce” the space charge problem’s complexity which ultimately limits the maximum attainable spatial resolution
- Most of the LSC studies use a simple 1D model based on impedance approximation
- Space charge problem is very similar to the well-known  $N$ -body problem in celestial mechanics
- **We used very effective algorithm for the gravitational  $N$ -body problem, so called “tree” or Barnes-Hut (BH) algorithm\***

Some conventional codes: ASTRA, SYNERGIA, TSTEP

\*J. Barnes and P. Hut, *Nature*, **324**, 446 (1986).

# Tree algorithm: in brief

- Scales as  $\mathcal{O}(N \log N)$ , where  $N$  is the number of macroparticles used to represent the beam
- Precision parameter corresponds to the “depth” of the tree
- Can be applied to many-body systems

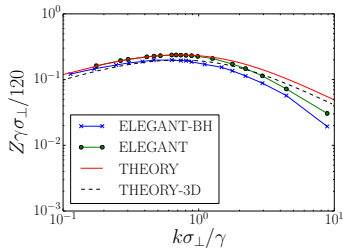
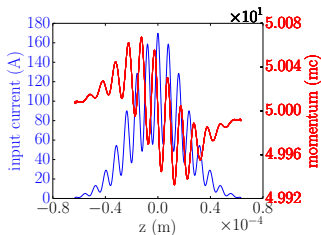


Images courtesy of J. Barnes



# Code validation

Let's consider initial bunch distribution with pre-modulated current profiles of the form  $f(\mathbf{r}) = T(x, y)L_z(z) [1 + m \cos kz]$



*On the left:* Initial density modulation resulted in energy modulation. *On the right:* The agreement between the BH algorithm and analytical impedance equation

$$Z(k) = -i \frac{Z_0}{\pi\gamma\sigma} \frac{\xi_{\sigma}}{4} e^{\xi_{\sigma}^2/2} \text{Ei}\left(-\frac{\xi_{\sigma}^2}{2}\right)$$

# Bunching factor and gain

To characterize the current (density) modulation one can introduce the bunching factor

$$b(\omega) = \frac{1}{N} \left| \sum_n \exp(-i\omega t_n) \right|$$

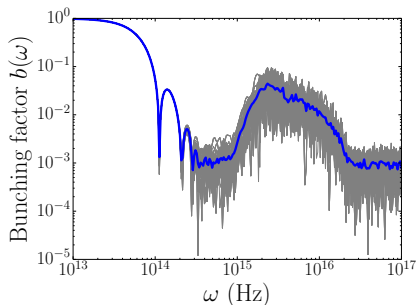
The broadband amplification process can be seen on the bunching factor curve as a broad peak. One can numerically compute the gain as:

$$G(\omega) = G_1 \times G_2 \times \dots \times G_n = \left| \frac{b_f(\omega_f)}{b_0(\omega_i)} \right|$$

\*JLAB-TN-14-016, Rui Li and C.-Y. Tsai

# Bunching factor (averaged)

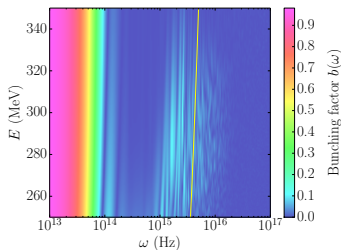
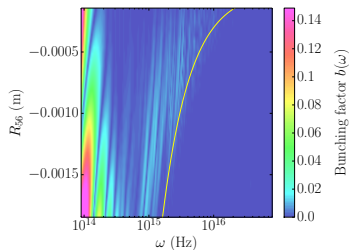
The LSC impedance results in selection of preferred frequency



100 realizations with 1M particles (gray traces) and corresponding average (blue trace)

# $b(\omega)$ as a function of E and chicane

*On the left:* Bunching factor for different values of the chicane long. dispersion  $R_{56}$



*On the right:* The change of the bunching factor vs energy of the bunch

*Yellow solid line is analytical prediction.*

More results: A. Halavanau and P. Piot, NIM A 2016 819 144-153.

# Desired bunch parameters

Parameter	Value	Units
Spotsize, $\sigma$	2.2 - 70.4	$\mu\text{m}$
Charge, $Q$	20.0	pC
Lorentz factor, $\gamma$	50 - 1000	–
Bunch duration, $\tau$	120	fs
Norm. transv. emittance, $\varepsilon_{x,y}$	$10^{-8}$	m
Momentum spread, $\sigma_\delta$	$10^{-4}$	–
Total LSCA length, $D$	28.0	m

# LSCA at FAST Summary

- Using a gridless code adapted from Astrophysics we have investigated effects in the LSC impedance and found that the one-dimensional often used LSC impedance model is a good approximation **(NEW)**
- Will not require much redesign of the lattice, can be compact (10-20 m), also will help to turn FAST injector into FEL
- We demonstrated that LSCA can produce femtosecond pulses of light in optical regime. Still needs to be pushed for the VUV regime **(NEW)**

# Final conclusions

- ① Existing analytical model of 1.3 GHz accelerating SRF cavity confirmed, backbone of ILC, LCLS-II
- ② Developed MLA based laser transverse shaping technique, significantly improved beam emittance
- ③ Generated CAM and flat beams at FAST, on way to ILC-type beams
- ④ Generated tunable bunch trains with MLA+EEX, many outcomes
- ⑤ Had a lot of fun



# Fermilab Accelerator PhD program



*Vita: 3 papers + 2 in progress*

- ① Simulation of a cascaded longitudinal space charge amplifier for coherent radiation generation, NIMA, 819, (2016) 144-153
- ② Analysis and Measurement of the Transfer Matrix of a 9-cell 1.3-GHz Superconducting Cavity, Phys. Rev. Accel. Beams 20, 4, 040102 (2017)
- ③ Spatial control of photoemitted electron beams using a microlens-array transverse-shaping technique, Phys. Rev. Accel. Beams 20, 103404 (2017)
- ④ Magnetized and flat beam experiment at FAST, IPAC2018, *paper in progress*
- ⑤ Simple technique for a tunable bunch train generation, IPAC18, *paper in progress*
- ⑥ **17 conference papers (first author)**



# Credits

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Thank you for your attention!